

§1 Introduction

Kraichnan, J. Fluid Mech. 5, 497-543 (1959) (DIA)

Kraichnan, Phys. Fluids 8, 575-598 (1965) (ALHDIA)

Kraichnan and Herring, J. Fluid Mech 88 355-367 (1978) (SBALHDIA)

Kaneda, J. Fluid Mech 107 131-145 (1981) (LRA).

Leslie, 'Developments in the theory of turbulence'
(Clarendon Press, Oxford, 1973)

Kaneda Fluid Dynamics Research 39 526-551 (2007) (Review)

後藤俊幸 「乱流理論の基礎」 (朝倉書店, 1992)

Incompressible Navier-Stokes equations

$$\frac{\partial u^\alpha}{\partial t} + u^\beta \frac{\partial}{\partial x^\beta} u^\alpha = - \frac{\partial p}{\partial x^\alpha} + \nu \frac{\partial^2 u^\alpha}{\partial x^\beta \partial x^\beta} + f^\alpha \quad (1)$$

$$\frac{\partial u^\alpha}{\partial x^\alpha} = 0 \quad (2)$$

$u^\alpha(x, t)$: velocity field

$p(x, t)$: pressure field

ν : kinematic viscosity

f^α : external force field

Fourier Transform

$$\begin{cases} u^\alpha(x) = \int_{\underline{k}} u_{\underline{k}}^\alpha e^{i\underline{k}\cdot\underline{x}} & (3) \end{cases}$$

$$\begin{cases} u_{\underline{k}}^\alpha = \frac{1}{(2\pi)^d} \int_{\underline{x}} u^\alpha(x) e^{-i\underline{k}\cdot\underline{x}} & (4) \end{cases}$$

Navier - Stokes equations in Fourier space

$$\begin{cases} \frac{\partial}{\partial t} u_{\underline{k}}^\alpha = \int_{\underline{p}, \underline{q}} \delta(\underline{k}-\underline{p}-\underline{q}) M_{\underline{k}}^{\alpha\beta\gamma} u_{\underline{p}}^\beta u_{\underline{q}}^\gamma \\ \quad - \nu k^2 u_{\underline{k}}^\alpha + D_{\underline{k}}^{\alpha\gamma} f_{\underline{k}}^\gamma & (5) \end{cases}$$

$$\underline{k}^\alpha u_{\underline{k}}^\alpha = 0 \quad (6)$$

$$M_{\underline{k}}^{\alpha\beta\gamma} := -\frac{i}{2} (k^\beta D_{\underline{k}}^{\alpha\gamma} + k^\gamma D_{\underline{k}}^{\alpha\beta}) \quad (7)$$

$$D_{\underline{k}}^{\alpha\beta} := \delta^{\alpha\beta} - \frac{k^\alpha k^\beta}{k^2} \quad (8)$$

reality condition

$$u_{-\underline{k}}^\alpha = u_{\underline{k}}^{\alpha*} \quad (9)$$

$$\frac{\partial}{\partial t} \frac{1}{2} |u_{\underline{k}}^\alpha|^2 = \frac{1}{2} \int_{\underline{p}, \underline{q}} \delta(\underline{k} - \underline{p} - \underline{q}) \left(\tilde{H}_{\underline{k}, \underline{p}, \underline{q}}^{\alpha\alpha} + \overset{\text{c.c.}}{\tilde{H}_{-\underline{k}, -\underline{p}, -\underline{q}}^{\alpha\alpha}} \right) - 2\nu k^2 |u_{\underline{k}}^\alpha|^2 + D_{\underline{k}}^{\alpha\gamma} f_{\underline{k}}^\gamma u_{\underline{k}}^\alpha + D_{\underline{k}}^{\alpha\gamma} u_{\underline{k}}^\alpha f_{\underline{k}}^\gamma \quad (10)$$

$$\tilde{H}_{\underline{k}, \underline{p}, \underline{q}}^{\alpha\beta} := M_{\underline{k}}^{\alpha\gamma\delta} u_{\underline{p}}^\gamma u_{\underline{q}}^\delta u_{-\underline{k}}^\beta \quad (11)$$

$$-\frac{1}{2} \left(\tilde{H}_{\underline{k}, \underline{p}, \underline{q}}^{\alpha\alpha} + \tilde{H}_{-\underline{k}, -\underline{p}, -\underline{q}}^{\alpha\alpha} \right) \quad (12)$$

↑
energy transfer rate $\underline{k} \rightarrow \underline{p}, \underline{q}$

$$\int_{\underline{k}, |\underline{k}| > K} (10) \quad \frac{\partial}{\partial t} \frac{1}{2} \int_{\underline{k}, |\underline{k}| > K} |u_{\underline{k}}^\alpha|^2 = \tilde{\Pi}(K) - \tilde{E}(K) + \tilde{F}(K) \quad (13)$$

• energy flux into $|\underline{k}| > K$

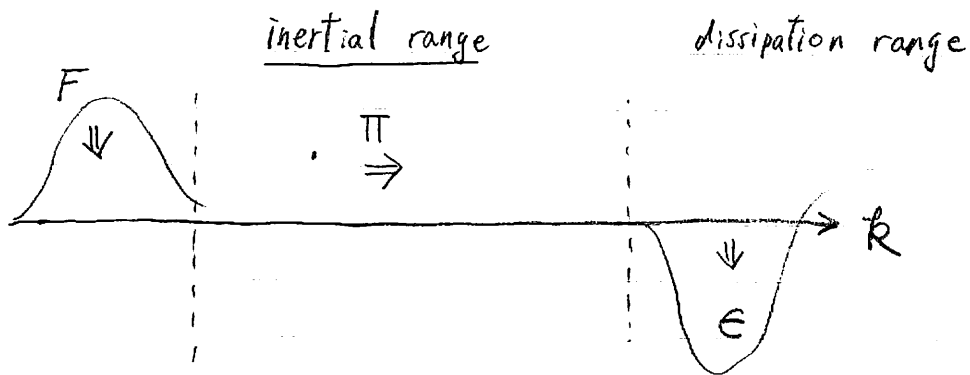
$$\tilde{\Pi}(K) := \frac{1}{2} \int_{\underline{k}, |\underline{k}| > K} \int_{\underline{p}, \underline{q}} \delta(\underline{k} - \underline{p} - \underline{q}) \left(\tilde{H}_{\underline{k}, \underline{p}, \underline{q}}^{\alpha\alpha} + \tilde{H}_{-\underline{k}, -\underline{p}, -\underline{q}}^{\alpha\alpha} \right) \quad (14)$$

• energy dissipation rate in $|\underline{k}| > K$

$$\tilde{E}(K) := \nu \int_{\underline{k}, |\underline{k}| > K} k^2 |u_{\underline{k}}^\alpha|^2 \quad (15)$$

• energy injection rate in $|\underline{k}| > K$

$$\tilde{F}(K) := \frac{1}{2} \int_{\underline{k}, |\underline{k}| > K} \left(D_{\underline{k}}^{\alpha\gamma} f_{\underline{k}}^\gamma u_{\underline{k}}^\alpha + D_{\underline{k}}^{\beta\gamma} u_{\underline{k}}^\alpha f_{-\underline{k}}^\gamma \right) \quad (16)$$



When k is in the inertial range.

$$\tilde{\epsilon}(k) \approx \tilde{\epsilon}(0) =: \tilde{\epsilon} \quad (17)$$

$$F(k) \approx 0 \quad (18)$$

(13), (17), (18)

$$\frac{\partial}{\partial t} \frac{1}{2} \int_{\underline{k}, |\underline{k}| > k} |u_{\underline{k}}^\alpha|^2 \approx \tilde{\pi}(k) - \tilde{\epsilon} \quad (19)$$

if

$$\frac{\partial}{\partial t} \frac{1}{2} \int_{\underline{k}, |\underline{k}| > k} |u_{\underline{k}}^\alpha|^2 \approx 0 \quad (20)$$

$$\tilde{\pi}(k) \approx \tilde{\epsilon} \quad (21)$$

Statistically averaged quantities

$$Q_{\underline{k}}^{\alpha\beta}(t) := \frac{1}{S(0)} \langle u_{\underline{k}}^{\alpha}(t) u_{-\underline{k}}^{\beta}(t) \rangle \quad (22) \quad \frac{1}{S(0)} \leftarrow \text{Homogeneity in space.}$$

$$\begin{aligned} H_{\underline{k}\underline{p}\underline{q}}^{\alpha\beta}(t) &:= \frac{1}{S(0)} \langle \tilde{H}_{\underline{k}\underline{p}\underline{q}}^{\alpha\beta}(t) \rangle \\ &= \frac{1}{S(0)} M_{\underline{k}}^{\alpha\gamma\delta} \langle u_{\underline{p}}^{\gamma}(t) u_{\underline{q}}^{\delta}(t) u_{-\underline{k}}^{\beta}(t) \rangle \quad (23) \end{aligned}$$

$$\frac{1}{S(0)} \langle (13) \rangle$$

$$\frac{1}{2} \frac{\partial}{\partial t} \int_{\underline{k}, |\underline{k}| > K} Q_{\underline{k}}^{\alpha\alpha}(t) = \Pi(K, t) - \epsilon(K, t) + F(K, t) \quad (24)$$

$$\begin{aligned} \Pi(K, t) &= \frac{1}{S(0)} \langle \tilde{\Pi}(K, t) \rangle \\ &= \frac{1}{2} \int_{\underline{k}, |\underline{k}| > K} \int_{\underline{p}, \underline{q}} S(\underline{k} - \underline{p} - \underline{q}) \left(H_{\underline{k}\underline{p}\underline{q}}^{\alpha\alpha} + H_{-\underline{k}\underline{p}-\underline{q}}^{\alpha\alpha} \right) \quad (25) \end{aligned}$$

$$\begin{aligned} \epsilon(K, t) &= \frac{1}{S(0)} \langle \tilde{\epsilon}(K, t) \rangle \\ &= \nu \int_{\underline{k}, |\underline{k}| > K} k^2 Q_{\underline{k}}^{\alpha\alpha}(t) \quad (26) \end{aligned}$$

$$F(K, t) = \frac{1}{S(0)} \langle F(K, t) \rangle \quad (27)$$

Need a model for $\Pi(K, t)$ or $H_{\underline{k}\underline{p}\underline{q}}^{\alpha\beta}(t)$.

Express $H_{\underline{k}\underline{p}\underline{q}}^{\alpha\beta}(t)$ in terms of $Q_{\underline{k}}^{\alpha\beta}(t)$

(and other quantities.)

§2 Derivation of Direct Interaction Approximation

Two-time correlation function

$$\langle u_{\underline{k}}^{\alpha}(t) u_{-\underline{k}'}^{\beta}(s) \rangle = Q_{\underline{k}}^{\alpha\beta}(t, s) \delta(\underline{k} - \underline{k}') \quad (1)$$

↖ Statistical homogeneity in space

Two-time response function

$$\langle \tilde{G}_{\underline{k}\underline{k}'}^{\alpha\beta}(t, s) \rangle = G_{\underline{k}}^{\alpha\beta}(t, s) \delta(\underline{k} - \underline{k}') \quad (2)$$

$$G_{\underline{k}}^{\alpha\beta} := \frac{\delta u_{\underline{k}}^{\alpha}(t)}{\delta f_{\underline{k}}^{\beta}(t)} \quad (3)$$

$$(1.5) \quad M \rightarrow \lambda M \quad (\lambda = 1)$$

$$\frac{\partial}{\partial t} u_{\underline{k}}^{\alpha} = \lambda \int_{\underline{p}, \underline{q}} \delta(\underline{k} - \underline{p} - \underline{q}) M_{\underline{k}\underline{p}\underline{q}}^{\alpha\beta\gamma} u_{\underline{p}}^{\beta} u_{\underline{q}}^{\gamma} - \nu k^2 u_{\underline{k}}^{\alpha} + D_{\underline{k}}^{\alpha\gamma} f_{\underline{k}}^{\gamma} \quad (4)$$

$$f = 0$$

$$\begin{aligned} \frac{\partial}{\partial t} Q_{\underline{k}}^{\alpha\beta}(t, t) &= \int_{\underline{p}, \underline{q}} \delta(\underline{k} - \underline{p} - \underline{q}) [H_{\underline{k}\underline{p}\underline{q}}^{\alpha\beta}(t) + H_{-\underline{k}-\underline{p}-\underline{q}}^{\beta\alpha}(t)] \\ &\quad - 2\nu k^2 Q_{\underline{k}}^{\alpha\beta}(t, t) \end{aligned} \quad (5)$$

$$\frac{\partial}{\partial t} Q_{\underline{k}}^{\alpha\beta}(t, s) = \int_{\underline{p}, \underline{q}} \delta(\underline{k} - \underline{p} - \underline{q}) I_{\underline{k}\underline{p}\underline{q}}^{\alpha\beta}(t, s) - \nu k^2 Q_{\underline{k}}^{\alpha\beta}(t, s) \quad (6)$$

$$\frac{\partial}{\partial t} G_{\underline{k}}^{\alpha\beta}(t, s) = \int_{\underline{p}, \underline{q}} \delta(\underline{k} - \underline{p} - \underline{q}) J_{\underline{k}\underline{p}\underline{q}}^{\alpha\beta}(t, s) - \nu k^2 G_{\underline{k}}^{\alpha\beta}(t, s) \quad (7)$$

$$-G_{\underline{k}}^{\alpha\beta}(t, t) = D_{\underline{k}}^{\alpha\beta} \quad (8)$$

$$H_{\underline{k} \underline{p} \underline{q}}^{\alpha \beta}(t) := \lambda \frac{1}{S(0)} M_{\underline{k}}^{\alpha \gamma \delta} \langle u_{\underline{p}}^{\gamma}(t) u_{\underline{q}}^{\delta}(t) u_{-\underline{k}}^{\beta}(t) \rangle \quad (9)$$

$$I_{\underline{k} \underline{p} \underline{q}}^{\alpha \beta}(t, s) := \lambda \frac{1}{S(0)} M_{\underline{k}}^{\alpha \gamma \delta} \langle u_{\underline{p}}^{\gamma}(t) u_{\underline{q}}^{\delta}(t) u_{-\underline{k}}^{\beta}(t) \rangle \quad (10)$$

$$J_{\underline{k} \underline{p} \underline{q}}^{\alpha \beta}(t, s) := \lambda \frac{2}{S(0)} M_{\underline{k}}^{\alpha \gamma \delta} \langle \tilde{G}_{\underline{p} \underline{k}}^{\gamma \beta}(t, s) u_{\underline{q}}^{\delta}(t) \rangle \quad (11)$$

Purpose approximate H, I, J in terms of Q, G .

Primitive expansion in λ

$$u_{\underline{k}}^{\alpha} = u_{\underline{k}}^{\alpha(0)} + \lambda u_{\underline{k}}^{\alpha(1)} + \lambda^2 u_{\underline{k}}^{\alpha(2)} + \dots \quad (12)$$

$$\tilde{G}_{\underline{k} \underline{k}'}^{\alpha \beta} = \tilde{G}_{\underline{k} \underline{k}'}^{\alpha \beta(0)} + \lambda \tilde{G}_{\underline{k} \underline{k}'}^{\alpha \beta(1)} + \lambda^2 \tilde{G}_{\underline{k} \underline{k}'}^{\alpha \beta(2)} + \dots \quad (13)$$

From (7), (8), (13)

$$G_{\underline{k} \underline{k}'}^{\alpha \beta(0)}(t, s) = \begin{cases} e^{-\nu k^2(t-s)} D_{\underline{k}}^{\alpha \beta} S(\underline{k} - \underline{k}') & (t \geq s) \\ 0 & (t < s) \end{cases} \quad (14)$$

From (4), (14)

$$u_{\underline{k}}^{\alpha}(t) = u_{\underline{k}}^{\alpha(0)}(t) + \lambda \int_{t_0}^t dt' \int_{\underline{k}'} \tilde{G}_{\underline{k} \underline{k}'}^{\alpha \beta(0)}(t, t') \int_{\underline{p}', \underline{q}'} S(\underline{k}' - \underline{p}' - \underline{q}') M_{\underline{k}'}^{\beta \gamma \delta} u_{\underline{p}'}^{\gamma}(t') u_{\underline{q}'}^{\delta}(t') \quad (15)$$

$$u_{\underline{k}}^{\alpha(0)}(t) = \int_{\underline{k}'} \bar{G}_{\underline{k} \underline{k}'}^{\alpha \beta(0)}(t, t_0) u_{\underline{k}'}^{\beta}(t_0) \quad (16)$$

$$\frac{\delta}{\delta f_{\underline{k}'}^\beta(s)} \quad (15) \quad (t_0 \rightarrow s)$$

$$\begin{aligned} \tilde{G}_{\underline{k}\underline{k}'}^{\alpha\beta}(t,s) &= \tilde{G}_{\underline{k}\underline{k}'}^{\alpha\beta}(t,s) \\ &+ 2\lambda \int_s^t dt' \int_{\underline{k}''} G_{\underline{k}\underline{k}''}^{\alpha\eta(0)}(t,t') \int_{\underline{p}''\underline{q}''} \delta(\underline{k}'' - \underline{p}'' - \underline{q}'') \\ &\quad \times M_{\underline{k}''}^{\eta\gamma\zeta} \tilde{G}_{\underline{p}''\underline{k}'}^{\gamma\beta}(t',s) u_{\underline{q}''}^\zeta(t') \end{aligned} \quad (17)$$

Assumption

Odd moments

$$\langle u_{\underline{k}}^{\alpha(0)} \rangle = 0, \quad \langle u_{\underline{k}}^{\alpha(0)} u_{\underline{p}}^{\beta(0)} u_{\underline{q}}^{\gamma(0)} \rangle = 0, \quad \dots \quad (18)$$

$$\begin{aligned} \langle u_{\underline{k}}^{\alpha(0)} u_{\underline{p}}^{\beta(0)} u_{\underline{q}}^{\gamma(0)} u_{\underline{r}}^{\zeta(0)} \rangle &= \langle u_{\underline{k}}^{\alpha(0)} u_{\underline{p}}^{\beta(0)} \rangle \langle u_{\underline{q}}^{\gamma(0)} u_{\underline{r}}^{\zeta(0)} \rangle \\ &+ \langle u_{\underline{k}}^{\alpha(0)} u_{\underline{q}}^{\gamma(0)} \rangle \langle u_{\underline{p}}^{\beta(0)} u_{\underline{r}}^{\zeta(0)} \rangle \\ &+ \langle u_{\underline{k}}^{\alpha(0)} u_{\underline{r}}^{\zeta(0)} \rangle \langle u_{\underline{p}}^{\beta(0)} u_{\underline{q}}^{\gamma(0)} \rangle \end{aligned} \quad (19)$$

Primitive expansion

$$Q_{\underline{k}}^{\alpha\beta}(t,s) = Q_{\underline{k}}^{\alpha\beta(0)}(t,s) + \lambda Q_{\underline{k}}^{\alpha\beta(1)}(t,s) [Q^{(0)}, G^{(0)}] + O(\lambda^2) \quad (20)$$

$$G_{\underline{k}}^{\alpha\beta}(t,s) = G_{\underline{k}}^{\alpha\beta(0)}(t,s) + \lambda G_{\underline{k}}^{\alpha\beta(1)}(t,s) [Q^{(0)}, G^{(0)}] + O(\lambda^2) \quad (21)$$

$$H_{\underline{k}\underline{l}\underline{m}}^{\alpha\beta}(t) = \lambda^2 H_{\underline{k}\underline{l}\underline{m}}^{\alpha\beta(2)}(t) [Q^{(0)}, G^{(0)}] + O(\lambda^3) \quad (22)$$

$$I_{\underline{k}\underline{l}\underline{m}}^{\alpha\beta}(t,s) = \lambda^2 I_{\underline{k}\underline{l}\underline{m}}^{\alpha\beta(2)}(t,s) [Q^{(0)}, G^{(0)}] + O(\lambda^3) \quad (23)$$

$$J_{\underline{k}\underline{l}\underline{m}}^{\alpha\beta}(t,s) = \lambda^2 J_{\underline{k}\underline{l}\underline{m}}^{\alpha\beta(2)}(t,s) [Q^{(0)}, G^{(0)}] + O(\lambda^3) \quad (24)$$

Inverse expansion

$$Q_{\underline{k}}^{\alpha\beta(0)}(t,s) = Q_{\underline{k}}^{\alpha\beta}(t,s) + \lambda Q_{\underline{k}}^{\alpha\beta[1]}(t,s) [Q, G] + O(\lambda^2) \quad (25)$$

$$G_{\underline{k}}^{\alpha\beta(0)}(t,s) = G_{\underline{k}}^{\alpha\beta}(t,s) + \lambda G_{\underline{k}}^{\alpha\beta[1]}(t,s) [Q, G] + O(\lambda^2) \quad (26)$$

Renormalized expansion

(22) \leftarrow (25), (26)

$$H_{\underline{k}\underline{l}\underline{q}}^{\alpha\beta}(t) = \lambda^2 H_{\underline{k}\underline{l}\underline{q}}^{\alpha\beta(2)}(t) [Q, G] + O(\lambda^3) \quad (27)$$

(23) \leftarrow (25), (26)

$$I_{\underline{k}\underline{l}\underline{q}}^{\alpha\beta}(t, s) = \lambda^2 I_{\underline{k}\underline{l}\underline{q}}^{\alpha\beta(2)}(t, s) [Q, G] + O(\lambda^3) \quad (28)$$

(24) \leftarrow (25), (26)

$$J_{\underline{k}\underline{l}\underline{q}}^{\alpha\beta}(t, s) = \lambda^2 J_{\underline{k}\underline{l}\underline{q}}^{\alpha\beta(2)}(t, s) [Q, G] + O(\lambda^3) \quad (29)$$

Truncation

Retain the non-trivial leading order terms.

i.e. $O(\lambda^2)$

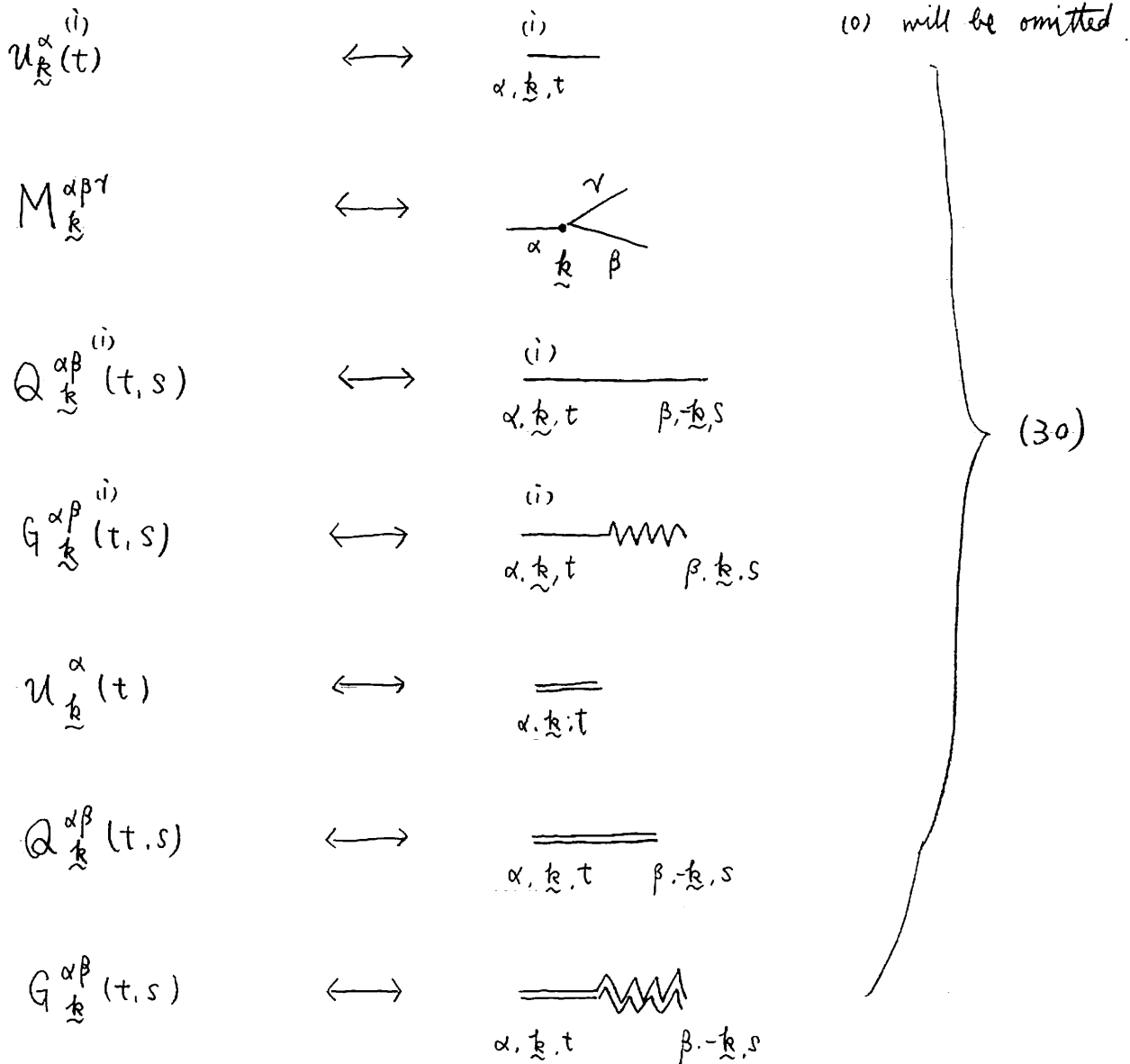
\Downarrow

Direct Interaction Approximation (DIA).

H, I, J are expressed in terms of Q, G .

Thus. Eqs. (5) - (8) are closed.

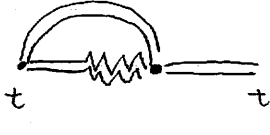

Diagram expression



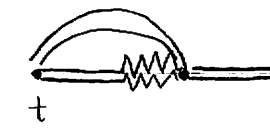
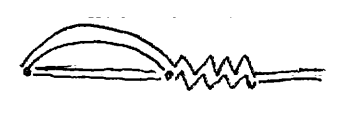
(30)

Renormalized expansion


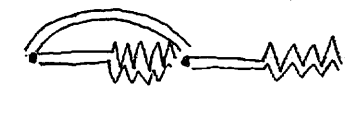
$$\left(\frac{\partial}{\partial t} + 2\nu k^2\right) \overline{\overline{t}}_t = \lambda^2 \left(4 \text{ [diagram 1]} + 2 \text{ [diagram 2]} + \text{c.c.} \right) + o(\lambda^3) \quad (42)$$

$$\left(\frac{\partial}{\partial t} + \nu k^2\right) \overline{\overline{t}}_s = \lambda^2 \left(4 \text{ [diagram 1]} + 2 \text{ [diagram 2]} \right) + o(\lambda^3) \quad (43)$$

$$\left(\frac{\partial}{\partial t} + \nu k^2\right) \overline{\overline{t}}_t \text{ [diagram 3]} = \lambda^2 4 \text{ [diagram 4]} + o(\lambda^3) \quad (44)$$

$$\begin{aligned}
 \left(\frac{\partial}{\partial t} + 2\nu k^2 \right) Q_{\underline{k}}^{\alpha\beta}(t) &= \lambda^2 \int_{\underline{p}, \underline{q}} \delta(\underline{k} - \underline{p} - \underline{q}) \int_{t_0}^t dt' \\
 &\times \left[4 M_{\underline{k}}^{\alpha\gamma\zeta} M_{\underline{p}}^{\eta\theta\kappa} G_{\underline{p}}^{\gamma\eta}(t, t') Q_{\underline{q}}^{\zeta\theta}(t, t') Q_{-\underline{k}}^{\beta\kappa}(t, t') \right. \\
 &\quad \left. - 2 M_{\underline{k}}^{\alpha\gamma\zeta} M_{\underline{k}}^{\kappa\eta\theta} Q_{\underline{p}}^{\gamma\eta}(t, t') Q_{\underline{q}}^{\zeta\theta}(t, t') G_{-\underline{k}}^{\beta\kappa}(t, t') \right] \\
 &+ (\underline{k} \leftrightarrow -\underline{k}, \alpha \leftrightarrow \beta) + O(\lambda^3) \quad (45)
 \end{aligned}$$

$$\begin{aligned}
 \left(\frac{\partial}{\partial t} + \nu k^2 \right) Q_{\underline{k}}^{\alpha\beta}(t, s) &= \lambda^2 \int_{\underline{p}, \underline{q}} \delta(\underline{k} - \underline{p} - \underline{q}) \int_{t_0}^t dt' \\
 &\times \left[4 M_{\underline{k}}^{\alpha\gamma\zeta} M_{\underline{p}}^{\eta\theta\kappa} G_{\underline{p}}^{\gamma\eta}(t, t') Q_{\underline{q}}^{\zeta\theta}(t, t') Q_{-\underline{k}}^{\beta\kappa}(s, t') \right. \\
 &\quad \left. - 2 M_{\underline{k}}^{\alpha\gamma\zeta} M_{\underline{k}}^{\kappa\eta\theta} Q_{\underline{p}}^{\gamma\eta}(t, t') Q_{\underline{q}}^{\zeta\theta}(t, t') G_{-\underline{k}}^{\beta\kappa}(s, t') \right] \\
 &+ O(\lambda^3) \quad (46)
 \end{aligned}$$

$$\begin{aligned}
 \left(\frac{\partial}{\partial t} + \nu k^2 \right) G_{\underline{k}}^{\alpha\beta}(t, s) &= \lambda^2 \int_{\underline{p}, \underline{q}} \delta(\underline{k} - \underline{p} - \underline{q}) \int_s^t dt' \\
 &\times \left[4 M_{\underline{k}}^{\alpha\gamma\zeta} M_{\underline{p}}^{\eta\theta\kappa} G_{\underline{p}}^{\gamma\eta}(t, t') Q_{\underline{q}}^{\zeta\theta}(t, t') G_{-\underline{k}}^{\kappa\beta}(t', s) \right] \\
 &+ O(\lambda^3) \quad (47)
 \end{aligned}$$

§3 Properties of DIA

✓ Energy conservation ($\nu = 0$)

✓ When $\nu = 0$, compatible with

Equipartition

$$Q_{\underline{k}}^{\alpha\beta}(t, t) = C D_{\underline{k}}^{\alpha\beta} \quad \text{--- (1)}$$

and Fluctuation-dissipation Relation

$$Q_{\underline{k}}^{\alpha\beta}(t, s) = G_{\underline{k}}(t, s) Q_{\underline{k}}^{\alpha\beta}(s, s) \quad \text{--- (2)}$$

where

$$G_{\underline{k}}^{\alpha\beta}(t, s) = G_{\underline{k}}(t, s) D_{\underline{k}}^{\alpha\beta} \quad \text{--- (3)}$$

✓ Realizability

i.e. compatible with a Langevin model.

✓ DIA can describe energy cascade

X Energy Spectrum $E(k) \propto U \overbrace{\epsilon^{1/2}}^{\langle u_i^2 \rangle} k^{-3/2} \quad \text{(4)}$

Inconsistent with the Kolmogorov spectrum

$$E(k) \propto \epsilon^{2/3} k^{-5/3} \quad \text{(5)}$$

Energy spectrum in inertial range.

Assume Isotropy

$$Q_{\underline{k}}^{\alpha\beta}(t,s) = \frac{1}{d-1} Q_k(t,s) D_{\underline{k}}^{\alpha\beta} \quad (6)$$

$$G_{\underline{k}}^{\alpha\beta}(t,s) = G_k(t,s) D_{\underline{k}}^{\alpha\beta} \quad (7)$$

It can be shown that $\int_{\underline{r}, \underline{z}}$ integrals in (2.46) and (2.47) diverge as $g \rightarrow 0$,

when

$$Q_k(t,t) \propto k^a \quad (8)$$

with

$$a < -d. \quad (9)$$

Then, (2.46), (2.47) reduce to

$$\left(\frac{\partial}{\partial t} + \nu k^2\right) G_k(t,s) = -C_d U^2 k^2 \int_s^t dt' G_k(t,t') G_k(t',s) \quad (10)$$

$$Q_k(t,s) = G_k(t,s) Q_k(s,s) \quad (11)$$

$C_d \dots \text{const.}$

$$U^2 = \frac{1}{d} \int_{\underline{g}} Q_{\underline{g}}(t,t) \quad (12)$$

\leftarrow Large scale quantity
 \uparrow
 outside inertial range

Similarity solution

$$G_k(t, s) = g(Uk(t-s)) \quad (13)$$

$$\frac{\partial}{\partial \tau} g(\tau) = -C_d \int_0^\tau d\tau' g(\tau-\tau') g(\tau') \quad (14)$$

$$g(0) = 1 \quad (15)$$

$$g(\tau) = \frac{J_1(2\sqrt{C_d}\tau)}{\tau} \quad (16)$$

J_1 : Bessel function of the first kind
of the order 1.

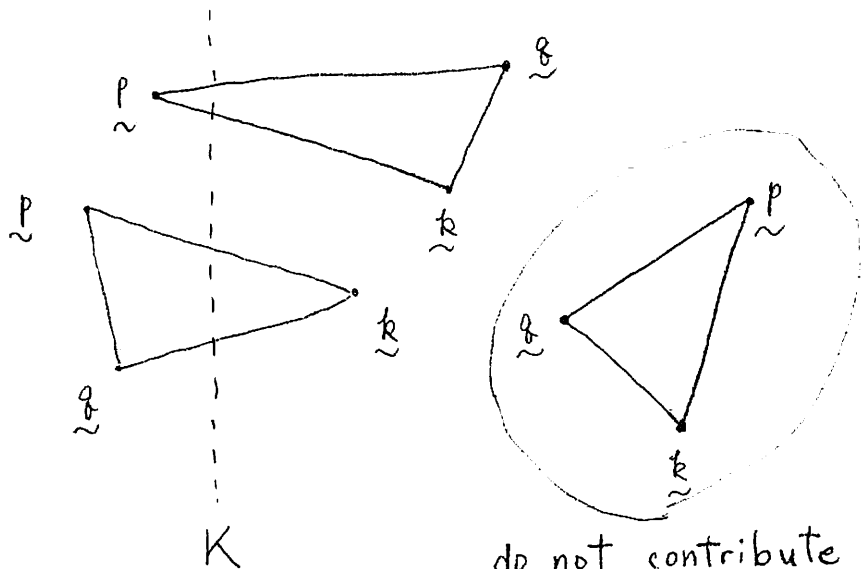
Energy Flux

(1.25)

$$\begin{aligned}
 \Pi(K, t) &= \frac{1}{2} \int_{\underline{k}, |\underline{k}| > K} \int_{\underline{p}, \underline{q}} \delta(\underline{k} - \underline{p} - \underline{q}) (H_{\underline{k}\underline{p}\underline{q}}^{\alpha\alpha} + H_{-\underline{k}-\underline{p}-\underline{q}}^{\alpha\alpha}) \\
 &= \frac{1}{2} \int_{\underline{k}, |\underline{k}| > K} \int_{\underline{p}, \underline{q}} \delta(\underline{k} - \underline{p} - \underline{q}) \int_{t_0}^t dt' \\
 &\quad \times [4 M_{\underline{k}}^{\alpha\gamma\zeta} M_{\underline{p}}^{\eta\theta\kappa} G_{\underline{p}}^{\gamma\eta}(t, t') Q_{\underline{q}}^{\zeta\theta}(t, t') Q_{-\underline{k}}^{\theta\kappa}(t, t') \\
 &\quad - 2 M_{\underline{k}}^{\alpha\gamma\zeta} M_{\underline{k}}^{\kappa\eta\theta} Q_{\underline{p}}^{\gamma\eta}(t, t') Q_{\underline{q}}^{\zeta\theta}(t, t') G_{-\underline{k}}^{\theta\kappa}(t, t')] \\
 &\quad + (\underline{k} \leftrightarrow -\underline{k}, \alpha \leftrightarrow \beta) \tag{17}
 \end{aligned}$$

Detailed energy balance

$$H_{\underline{k}\underline{p}\underline{q}}^{\alpha\alpha} + H_{-\underline{p}\underline{q}\underline{k}}^{\alpha\alpha} + H_{-\underline{q}\underline{k}\underline{p}}^{\alpha\alpha} = 0 \tag{18}$$



do not contribute
to the integral in (17)

$$\Pi(K, t) = - \int_{\underline{k}, |\underline{k}| > K} \int_{\underline{p}, |\underline{p}| > K} \int_{\underline{q}, |\underline{q}| < |\underline{p}|} \delta(\underline{k} - \underline{p} - \underline{q}) \times (H_{\underline{k}\underline{p}\underline{q}}^{\alpha\alpha} + H_{-\underline{k}-\underline{p}-\underline{q}}^{\alpha\alpha}) \quad (19)$$

• Assume the similarity form in the inertial range

$$G_k(t, s) = g(U k(t-s)) \quad (k_0 \ll k \ll k_1) \quad (20)$$

$$Q_k(t, s) = G_k(t, s) Q_k(s, s) \quad (\quad " \quad) \quad (21)$$

$$Q_k(t, t) \propto U^b \epsilon^c k^a \quad (\quad " \quad) \quad (22)$$

• Assume the $\underline{k}, \underline{p}, \underline{q}$ integral in (19) converge as

$$k_0 \rightarrow 0 \quad \text{and} \quad k_1 \rightarrow \infty.$$

• Assume the constant energy flux

$$\Pi(K) = \epsilon \quad (k_0 \ll K \ll k_1) \quad (23)$$

Dimensional analysis

$$\begin{aligned} [E] &= [\Pi(K)] = \left[\int_{\underline{k}, |\underline{k}| > K} \int_{\underline{p}, \underline{q}} \delta(\underline{k} - \underline{p} - \underline{q}) M M \int dt G Q Q \right] \\ &= [K^d K^d K^d K^{-d} K K (K^{-1} U^{-1}) (U^b \epsilon^c k^a)^2] \\ &= [K^{2d+1+2a} U^{2b-1} \epsilon^{2c}] \quad (24) \end{aligned}$$

$$a = -d - \frac{1}{2}, \quad b = \frac{1}{2}, \quad c = \frac{1}{2} \quad (25)$$

$$Q_{\mathbf{k}}(t, t) \propto U \epsilon^{\frac{1}{2}} k^{-d - \frac{1}{2}} \quad (26)$$

Energy Spectrum

$$\begin{aligned} E(\mathbf{k}, t) &:= \frac{1}{2} \int_{\underline{\mathbf{k}'}} \delta(|\underline{\mathbf{k}'} - \mathbf{k}|) Q_{\underline{\mathbf{k}'}}^{\alpha\alpha}(t, t) \\ &= \frac{1}{2} \int_{\underline{\mathbf{k}'}} \delta(|\underline{\mathbf{k}'} - \mathbf{k}|) Q_{\mathbf{k}}(t, t) \\ &= \frac{1}{2} A_d k^{d-1} Q_{\mathbf{k}}(t, t) \end{aligned} \quad (27)$$

A_d : surface area of d -dim. unit sphere.

$$E(\mathbf{k}) \propto U^{\frac{1}{2}} \epsilon^{\frac{1}{2}} k^{-\frac{3}{2}} \quad (28)$$

Inconsistent with the Kolmogorov spectrum

$$E(\mathbf{k}) \propto \epsilon^{\frac{2}{3}} k^{-\frac{5}{3}} \quad (29)$$

This is due to including the sweeping effect which should not be included in the estimate of the energy transfer.

§ 4 Lagrangian renormalized approximation

Lagrangian variables

$$X(\underline{x}, s | t)$$

s : labeling time, t : measuring time

The quantity of X at time t of the fluid element which was at position \underline{x} at time s .

position function

$$\psi(\underline{x}', t; \underline{x}, s) := \delta(\underline{x}' - \underline{y}(\underline{x}, s | t)) \quad - (1)$$

$\underline{y}(\underline{x}, s | t)$: position at time t of the fluid element which was at position \underline{x} at time s .

$$\frac{\partial}{\partial t} \psi(\underline{x}', t; \underline{x}, s) = -u^\alpha(\underline{x}', t) \frac{\partial}{\partial x'^\alpha} \psi(\underline{x}', t; \underline{x}, s) \quad - (2)$$

$$X(\underline{x}, s | t) = \int d\underline{x}' X(\underline{x}', t) \psi(\underline{x}', t; \underline{x}, s)$$

— (3)

Two-point Lagrangian correlation function $Q^{(L)}$

Two-point Lagrangian response function $G^{(L)}$

$$Q^{\alpha\beta (L)}(\underline{x}, t; \underline{x}', t') := \begin{cases} \langle [Pu]^\alpha(\underline{x}, t | t) u^\beta(\underline{x}', t') \rangle & (t \geq t') \\ \langle u^\alpha(\underline{x}, t) [Pu]^\beta(\underline{x}', t | t') \rangle & (t < t') \end{cases} \quad (4)$$

$$G^{\alpha\beta (L)}(\underline{x}, t; \underline{x}', t') := \left\langle \frac{\delta [Pu]^\alpha(\underline{x}', t' | t)}{\delta f^\beta(\underline{x}', t')} \right\rangle \quad (5)$$

P : projection onto solenoidal component

In Fourier space

$$\begin{aligned}
 Q_{\underline{k}}^{\alpha\beta(L)}(t, t') &:= \frac{1}{(2\pi)^3} \int d(\underline{x}-\underline{x}') e^{-i\underline{k}(\underline{x}-\underline{x}')} Q^{\alpha\beta(L)}(\underline{x}, t; \underline{x}', t') \\
 &= \begin{cases} \frac{1}{S(\underline{e})} \langle D_{\underline{k}}^{\alpha\gamma} u_{\underline{k}}^{\gamma}(t|t') u_{-\underline{k}}^{\beta}(t') \rangle & (t \geq t') \\ \frac{1}{S(\underline{e})} \langle u_{\underline{k}}^{\alpha}(t) D_{\underline{k}}^{\beta\gamma} u_{-\underline{k}}^{\gamma}(t|t') \rangle & (t < t') \end{cases} \quad (6)
 \end{aligned}$$

$$\begin{aligned}
 G_{\underline{k}}^{\alpha\beta(L)}(t, t') &:= \int d(\underline{x}-\underline{x}') e^{-i\underline{k}(\underline{x}-\underline{x}')} G^{\alpha\beta(L)}(\underline{x}, t; \underline{x}', t') \\
 &= \frac{1}{S(\underline{e})} \langle \tilde{G}_{\underline{k}\underline{k}}^{\alpha\beta(L)}(t, t') \rangle \quad (t \geq t') \quad (7)
 \end{aligned}$$

$$\tilde{G}_{\underline{k}\underline{k}'}^{\alpha\beta(L)}(t, t') = \frac{\int D_{\underline{k}}^{\alpha\gamma} u_{\underline{k}}^{\gamma}(t)}{\int f_{\underline{k}'}^{\beta}(t')} \quad (8)$$

$$\psi_{\underline{k}, \underline{k}'}(t, t') := \frac{1}{(2\pi)^3} \int d\underline{x} \int d\underline{x}' e^{i\underline{k}\underline{x}'} e^{i\underline{k}'\underline{x}} \psi(\underline{x}, t; \underline{x}', t') \quad (9)$$

$$\psi_{\underline{k}}(t, t') = \frac{1}{S(\underline{e})} \langle \psi_{\underline{k}, -\underline{k}}(t, t') \rangle \quad (10)$$

$$\frac{\partial}{\partial t} Q_{\underline{k}}^{\alpha\beta}(t, t) = \int_{\underline{p}, \underline{q}} \delta(\underline{k} - \underline{p} - \underline{q}) \left[H_{\underline{k}, \underline{p}, \underline{q}}^{\alpha\beta}(t) + H_{-\underline{k}, -\underline{p}, -\underline{q}}^{\beta\alpha}(t) \right] - 2\nu \underline{k}^2 Q_{\underline{k}}^{\alpha\beta}(t, t) \quad (10)$$

$$\frac{\partial}{\partial t} Q_{\underline{k}}^{\alpha\beta(L)}(t, s) = \int_{\underline{p}, \underline{q}} \delta(\underline{k} - \underline{p} - \underline{q}) J_{\underline{k}, \underline{p}, \underline{q}}^{\alpha\beta(L)}(t, s) - \nu K_{\underline{k}}^{\alpha\beta}(t, s) \quad (11)$$

$$\frac{\partial}{\partial t} G_{\underline{k}}^{\alpha\beta(L)}(t, s) = \int_{\underline{p}, \underline{q}} \delta(\underline{k} - \underline{p} - \underline{q}) J_{\underline{k}, \underline{p}, \underline{q}}^{\alpha\beta(L)}(t, s) - \nu L_{\underline{k}}^{\alpha\beta}(t, s) \quad (12)$$

Primitive expansion

$$Q_{\underline{k}}^{\alpha\beta(L)}(t, s) = Q_{\underline{k}}^{\alpha\beta(0)}(t, s) + o(\lambda) \quad (13)$$

$$G_{\underline{k}}^{\alpha\beta(L)}(t, s) = G_{\underline{k}}^{\alpha\beta(0)}(t, s) + o(\lambda) \quad (14)$$

$$\psi_{\underline{k}}^{\alpha\beta(L)}(t, s) = \psi_{\underline{k}}^{\alpha\beta(0)}(t, s) + o(\lambda) \quad (15)$$

$$H_{\underline{k}, \underline{p}, \underline{q}}^{\alpha\beta}(t) = \lambda^2 H_{\underline{k}, \underline{p}, \underline{q}}^{\alpha\beta(2)}(t) [Q^{(0)}, G^{(0)}] + o(\lambda^3) \quad (16)$$

$$I_{\underline{k}, \underline{p}, \underline{q}}^{\alpha\beta(L)}(t) = \lambda^2 I_{\underline{k}, \underline{p}, \underline{q}}^{\alpha\beta(2)}(t) [Q^{(0)}, G^{(0)}, \psi^{(0)}] + o(\lambda^3) \quad (17)$$

$$J_{\underline{k}, \underline{p}, \underline{q}}^{\alpha\beta(L)}(t) = \lambda^2 J_{\underline{k}, \underline{p}, \underline{q}}^{\alpha\beta(2)}(t) [Q^{(0)}, G^{(0)}, \psi^{(0)}] + o(\lambda^3) \quad (18)$$

Inverse expansion

$$Q_{\underline{k}}^{\alpha\beta (0)}(t, s) = Q_{\underline{k}}^{\alpha\beta (L)}(t, s) + O(\lambda) \quad (19)$$

$$G_{\underline{k}}^{\alpha\beta (0)}(t, s) = G_{\underline{k}}^{\alpha\beta (L)}(t, s) + O(\lambda) \quad (20)$$

$$\left(\Psi_{\underline{k}}^{(0)}(t, s) = \Psi_{\underline{k}}(t, s) + O(\lambda) \right) \quad (21)$$

Renormalized expansion

$$H_{\underline{k}L\underline{p}\underline{q}}^{\alpha\beta}(t) = \lambda^2 H_{\underline{k}L\underline{p}\underline{q}}^{\alpha\beta}(t, s) [Q^{(L)}, G^{(L)}] + O(\lambda^3) \quad (22)$$

$$I_{\underline{k}L\underline{p}\underline{q}}^{\alpha\beta (L)}(t) = \lambda^2 I_{\underline{k}L\underline{p}\underline{q}}^{\alpha\beta}(t, s) [Q^{(L)}, G^{(L)}, \Psi^{(0)}] + O(\lambda^3) \quad (23)$$

$$J_{\underline{k}L\underline{p}\underline{q}}^{\alpha\beta}(t) = \lambda^2 J_{\underline{k}L\underline{p}\underline{q}}^{\alpha\beta}(t, s) [Q^{(L)}, G^{(L)}, \Psi^{(0)}] + O(\lambda^3) \quad (24)$$

Truncation

Retain the non-trivial leading order terms
in (22) (23) (24).

$$(10) (11) (12) \leftarrow (22) (23) (24)$$

↓

LRA equation.

§5 Properties of LRA

✓ Energy conservation ($\nu = 0$)

✓ When $\nu = 0$, compatible with equipartition (3.1) and

FD relation (3.2).

• Realizability \rightarrow not shown.

✓ Energy spectrum

$$E(k) = K_0 \epsilon^{2/3} k^{-5/3} \quad (1)$$

$$K_0 \doteq 1.72 \quad (2)$$

Kolmogorov spectrum.

This is because

$\int_{p,q} \dots$ integrals in (12)

$$\left. \begin{array}{l} \text{converge} \\ \text{diverge} \end{array} \right\} \text{ for } \left. \begin{array}{l} a > -d-2 \\ a < -d-2 \end{array} \right\} \quad (3)$$

as $q \rightarrow 0$ when $Q_{\mathbf{k}}(t, \tau) \propto k^a$

Similarity form in the inertial range

$$G_k^{(L)}(t,s) = g(\epsilon^e k^c (t-s)) \quad (4)$$

$$Q_k^{(L)}(t,s) = G_k^{(L)}(t,s) Q_k(s,s) \quad (5)$$

$$Q_k(t,t) \propto \epsilon^b k^a \quad (6)$$

$$\begin{aligned} [E] &= [\Pi(K)] = \left[\int_k |k| > k \int_{k, \underline{q}} \delta(k-p-q) M M \int dt G Q Q \right] \\ &= [K^d K^d K^d K^{-d} K K (\epsilon^{-e} K^{-c}) (\epsilon^b K^a)^2] \\ &= [K^{2d+2-c+2a} \epsilon^{-e+2b}] \quad (7) \end{aligned}$$

$$2d+2-c+2a = 0 \quad (8)$$

$$-e+2b = 1 \quad (9)$$

$$[E] = \left[\frac{L^2}{T^3} \right] \quad [k] = \left[\frac{1}{L} \right]$$

$$\left[\frac{1}{T} \right] = \left[\epsilon^{1/3} k^{2/3} \right]$$

$$e = 1/3 \quad c = 2/3 \quad (10)$$

$$a = -d - \frac{2}{3} \quad , \quad b = \frac{2}{3} \quad (11)$$

$$Q_k(t,t) \propto \epsilon^{2/3} k^{-d-2/3} \quad (12)$$

$$E(k,t) \propto \epsilon^{2/3} k^{-5/3} \quad (13)$$

§ 6 Wave turbulence

$$\begin{aligned}
 \frac{\partial}{\partial t} u_{\underline{k}}^{\alpha} = & \lambda \int_{\underline{p}, \underline{q}} S(\underline{k} - \underline{p} - \underline{q}) M_{\underline{k}}^{\alpha\beta\gamma} u_{\underline{p}}^{\beta} u_{\underline{q}}^{\gamma} \\
 & - \nu k^2 u_{\underline{k}}^{\alpha} + D_{\underline{k}}^{\alpha\gamma} f_{\underline{k}}^{\gamma} \\
 & + \underline{\underline{L}}_{\underline{k}}^{\alpha\beta} u_{\underline{k}}^{\beta}
 \end{aligned} \tag{1}$$

Linearized equation

$$\frac{\partial}{\partial t} u_{\underline{k}}^{\alpha} = L_{\underline{k}}^{\alpha\beta} u_{\underline{k}}^{\beta} \tag{2}$$

$$u_{\underline{k}}(t) \sim e^{i\omega_{\underline{k}} t} u_{\underline{k}} \tag{3}$$

↑

Wave.

When L is large compared to the nonlinear term, λ -expansion can be justified

$$v=0$$

$$G_{\underline{k}}^{(0)}(t, t') \sim e^{i\omega_{\underline{k}}(t-t')} \quad (4)$$

$$Q_{\underline{k}}(t, t') \sim G_{\underline{k}}^{(0)}(t, t') Q_{\underline{k}}(t', t') \quad (5)$$

$$\begin{aligned} \frac{\partial}{\partial t} Q_{\underline{k}}(t, t) &= \int_{\underline{p}, \underline{q}} \delta(\underline{k}-\underline{p}-\underline{q}) \int_{t_0}^t dt' M M G^{(0)}(t, t') Q(t, t') Q(t, t') \\ &= \int_{\underline{p}, \underline{q}} \delta(\underline{k}-\underline{p}-\underline{q}) \int_{t_0}^t dt' M M G^{(0)}(t, t') G^{(0)}(t, t') G^{(0)}(t, t') \\ &\quad \times \underbrace{Q(t', t') Q(t', t')}_{\downarrow} \quad (6) \\ &\quad \text{slowly varying} \rightarrow Q(t, t) Q(t, t) \end{aligned}$$

$$Q_{\underline{k}} = [\epsilon^b \omega_{\underline{k}}^c k^a] \quad (7)$$

$$\begin{aligned} [\epsilon] = [\Pi(K)] &= \left[\int_{\underline{k} \quad |\underline{k}| > K} \int_{\underline{p}, \underline{q}} \delta(\underline{k}-\underline{p}-\underline{q}) M M \int dt G^{(0)} Q Q \right] \\ &= [K^d K^d K^d K^{-d} K^2 \omega_{\underline{k}}^{-1} (\epsilon^b \omega_{\underline{k}}^c k^a)^2] \quad (8) \end{aligned}$$

$$\left\{ \begin{aligned} 2d+2+2a &= 0 \\ 2b &= 1 \\ -1+2c &= 0 \end{aligned} \right. \quad (9)$$

$$a = -d-1, \quad b = \frac{1}{2}, \quad c = \frac{1}{2} \quad (10)$$

$$Q_k \propto \epsilon^{1/2} \omega_k^{1/2} k^{-d-1} \quad (11)$$

$$E(k) \propto k^{d-1} Q_k \propto \epsilon^{1/2} \omega_k^{1/2} k^{-2} \quad (12)$$